

Target-In/Target-Out running –Statistical Considerations

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Let $N_{\text{beam}}^{\text{in,out}}$ be the number of beam particles sampled with the target in and out.

Let $N_{\text{ev}}^{\text{in,out}}$ be the number of events recorded with the target in and out.

Let $N_{\text{good}}^{\text{in,out}}$ be the number of good events that survive after all the cuts for the in,out samples.

Both target-in and target-out good events have had all the offline cuts placed on them including vertex constraints.

Let λ be the target interaction length. λ is approximately 0.01.

Then let $N_{\text{ev}}^{\text{in,out}} = N_{\text{beam}}^{\text{in,out}} f^{\text{in,out}}$, where $f^{\text{in,out}}$ are the fractions of beam particles that get recorded on to disk. These numbers should be similar in magnitude.

Let $N_{\text{good}}^{\text{in,out}} = g^{\text{in,out}} N_{\text{ev}}^{\text{in,out}}$.

Since we are dead time limited, the running time in both target in and out cases is dictated by the number of events written to disk.

Let $N_{\text{ev}}^{\text{tot}} = N_{\text{ev}}^{\text{in}} + N_{\text{ev}}^{\text{out}}$. The plan is to optimize the fraction r of target out running for a fixed number of events written to disk.

$$N_{\text{ev}}^{\text{out}} = r N_{\text{ev}}^{\text{tot}}; \text{ and } N_{\text{ev}}^{\text{in}} = (1-r) N_{\text{ev}}^{\text{tot}}$$

We are trying to determine the cross section which is proportional to the interaction length λ where,

$$I = f^{\text{in}} g^{\text{in}} - f^{\text{out}} g^{\text{out}}$$

$$S_I^2 = (f^{\text{in}})^2 S_{g^{\text{in}}}^2 + (f^{\text{out}})^2 S_{g^{\text{out}}}^2 \text{ where } \sigma^2 \text{ denotes variance.}$$

$$g^{in} = \frac{N_{good}^{in}}{N_{ev}^{in}}$$

$$\mathbf{s}_{gin}^2 = \frac{\mathbf{s}_{N_{good}^{in}}^2}{N_{ev}^{in^2}} = \frac{N_{good}^{in}}{N_{ev}^{in^2}} = \frac{g^{in}}{N_{ev}^{in}}$$

$$\mathbf{s}_{gout}^2 = \frac{\mathbf{s}_{N_{good}^{out}}^2}{N_{ev}^{out^2}} = \frac{N_{good}^{out}}{N_{ev}^{out^2}} = \frac{g^{out}}{N_{ev}^{out}}$$

This leads to

$$\mathbf{s}_I^2 = \frac{1}{N_{ev}^{tot}} \left(\frac{f^{in^2} g^{in}}{1-r} + \frac{f^{out^2} g^{out}}{r} \right)$$

We need to minimize the error on the cross section with respect to r.

$$\frac{d\mathbf{s}_I^2}{dr} = \frac{1}{N_{ev}^{tot}} \left(\frac{f^{in^2} g^{in}}{(1-r)^2} - \frac{f^{out^2} g^{out}}{r^2} \right) = 0$$

This leads to

$$\frac{r^2}{(1-r)^2} = \frac{f^{out^2} g^{out}}{f^{in^2} g^{in}} \equiv k^2$$

leading to

$$r = \frac{k}{1+k}$$

$$k = \frac{f^{out}}{f^{in}} \sqrt{\frac{g^{out}}{g^{in}}}$$

Figure 1 shows the quantity r as a function of k^2 .

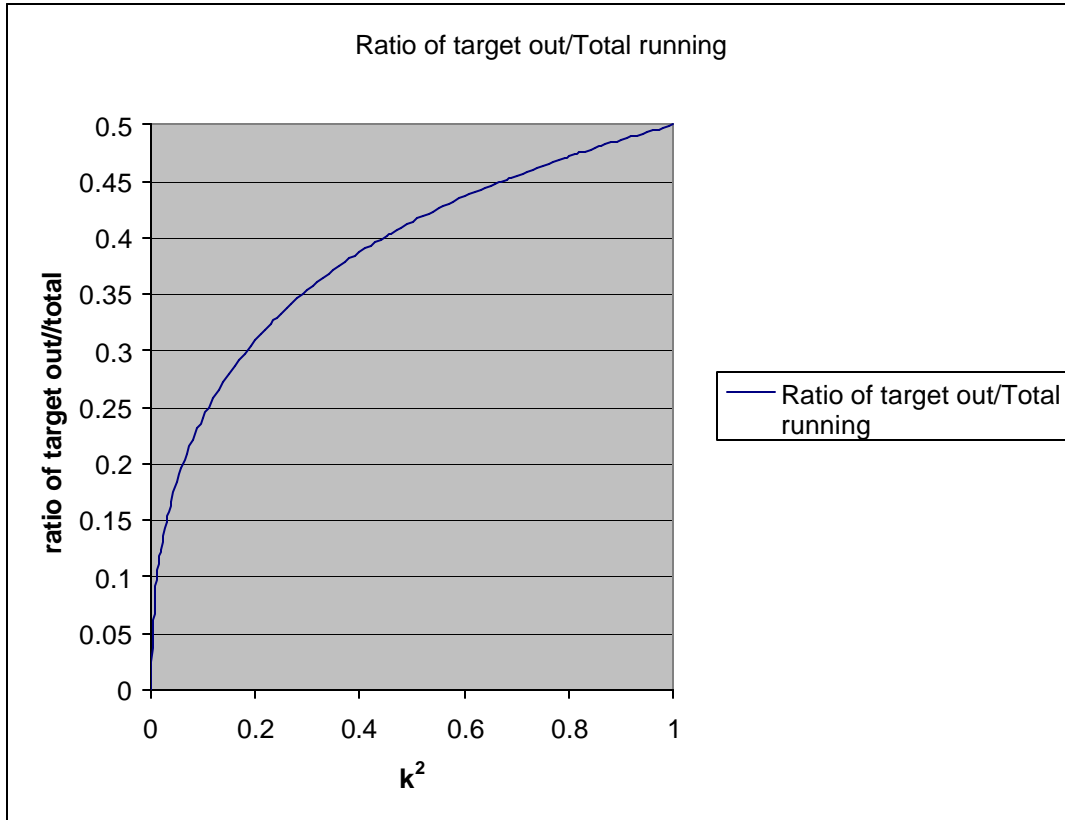


Figure 1 Optimal target-out running as a function of the ratio k^2 .

Notice the square root factor actually works against us, since g^{out}/g^{in} is less than unity.

With the optimum value of r , it is easy to show that

$$s_I^2 = \frac{f^{in^2} g^{in}}{N_{ev}^{tot} (1-r)} (1+k) = \frac{f^{in^2} g^{in}}{N_{ev}^{in}} (1+k)$$

i.e $1+k$ is the factor by which the purely target-in variance $\frac{f^{in^2} g^{in}}{N_{ev}^{in}}$ is inflated due to optimal target-out subtraction.